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A fractal set from the binary reflected Gray code

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Abstract

The permutation associated with the decimal expression of the binary reflected Gray code with N bits is considered. Its cycle structure is studied. Considered as a set of points, its self-similarity is pointed out. As a fractal, it is shown to be the attractor of an IFS. For large values of N the set is examined from the point of view of time series analysis.

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1. Introduction

In the most general setting one can define a Gray code as a listing of *words*, formed with N letters from an *alphabet* with k symbols, ordered in such a way that any two successive entries differ in just one symbol. These codes are named after Frank Gray, an engineer at Bell Laboratories who employed it in a patent in 1953 [10], although the use of the code has been traced back at least to 1878 when the French engineer Jean Maurice Emile Baudot (after whom the name baud for the signalling rate measure was chosen) developed a telegraph code on these grounds which readily prevailed in France.

We are interested in the case where numerical symbols are involved. We can still distinguish different types of codes according to the number system used. The most commonly employed bases are of course 2 and 10 giving rise to the *binary* and *decimal* Gray codes, respectively. Originally, the characteristic feature of these codes was important in mechanical devices such as counters, for no more than one piece at a time needs to act in every counting operation. Needless to say, this feature decreases the possibility of occurrence of errors albeit these may still be large. Although nowadays perhaps old fashioned, a very illustrative and familiar situation comes from the mechanical odometer of a car. When it reads 9999 km, five inner wheels have to rotate to show 10 000 km. Gray code is designed to avoid such *multiple* changes.

In a wide variety of types Gray codes have been and still are applied in a large range of extremely distinct fields. We leave apart their use as puzzle solvers [7] (e.g., the Towers of Hanoi or the Chinese rings) or in some even more bizarre application like campanology [26]. As a sample of more canonical fields where Gray codes are relevant we can cite as mere examples: database searching [17], design of charged particle detectors [19], space-filling Hilbert curves [8], genetic algorithms [12] and their use in parameter optimization problems [21], neural networks [16], computation of unstable periodic orbits in chaotic attractors [4] and, very recently, quantum computing [25]. Even a little surfing on the web will convince the reader of their importance in commercial matters related to A–D converters or design of communication codes.

In our case, it was using the native Gray code to generate Walsh functions that we were led to the considerations presented in this paper. We take a Gray code as the playground in which we build and represent some construction that is then analysed by techniques familiar in other fields like dynamical systems or time series analysis but whose application is novel, so we think, in this context. The construction we refer to is a permutation, P_G , of natural numbers which in fact constitutes a known sequence, listed as A003188 in OEIS [24]. Our approach is basically descriptive and is not directly related to the interesting topic of the mutual relationship between Gray codes and permutations which has been extensively considered in the literature [15]. In section 2, the specific type of Gray code which we will focus our attention on is briefly reviewed, namely the so-called *binary reflected Gray code* (BRGC for short). Then we build what we call Gray curve which we analyse from a fractal geometry point of view in section 3. Statistical considerations borrowed from modern time series analysis are applied in section 4. For easier reading we have postponed to section 5 the discussion of the technical details of the structure of the permutation of 2^N elements associated with the native Gray code. In the last section we add some final comments.

2. The binary reflected Gray code

The general definition of Gray code recalled in the introduction corresponds to what is known in the literature as a (N, k) -Gray code and can in fact be fulfilled by very different orderings. Consequently, one can find in the literature a considerable variety of members in the family of Gray codes and generalizations [22]. We are only interested in the length- N *binary* subclass which includes all 2^N words of N bits ordered in such a way that any two consecutive words differ in just 1 bit. We will even restrict our attention to the binary reflected Gray codes which are formally defined in a recursive way [14].

Let Γ_0 denote the empty string and Γ_N be the BRGC sequence of N -bit strings, then

$$\Gamma_{N+1} = 0\Gamma_N, 1\Gamma_N^R \quad (1)$$

where $0\Gamma_N$ denotes the sequence Γ_N with 0 prefixed to each string and $1\Gamma_N^R$ denotes the sequence Γ_N in reverse order with 1 prefixed to each string.

Perhaps the easiest way to seize the idea that underlies the reflected binary Gray code is through an example. Table 1 shows the first 16 decimal integers $0, 1, \dots, 15$ and their corresponding binary Gray code with $N = 4$. The obvious first column lists the natural decimal ordering we are familiar with. The second column, which we shall not use explicitly in the following, is included just to explicitly appreciate the differences between standard base 2 and Gray code orderings. The latter is shown in the fourth column. A first look at it reveals the aforementioned defining feature: adjacent entries differ in one, and only one, bit. The third column is incorporated to illustrate more clearly the building procedure from equation (1).

Table 1. Native Gray code construction for the first 2^4 natural numbers and its associated permutation P_G .

Decimal	Binary	Building procedure	Gray	P_G
0	0000	0	0000	0
1	0001	1	0001	1
2	0010	11	0011	3
3	0011	10	0010	2
4	0100	110	0110	6
5	0101	111	0111	7
6	0110	101	0101	5
7	0111	100	0100	4
8	1000	1100	1100	12
9	1001	1101	1101	13
10	1010	1111	1111	15
11	1011	1110	1110	14
12	1100	1010	1010	10
13	1101	1011	1011	11
14	1110	1001	1001	9
15	1111	1000	1000	8

The fifth column, which is central to our discussion, displays the decimal version of the previous column.

We have identified four blocks of rows in table 1 which will help in the explanation that follows, referred to the third column. The first block with two rows can be considered the seed for the Gray code, it represents Γ_1 from equation (1). The next block includes the previous one together with the added next two rows. When observing only the right-most bit in these latter one sees simply the previous two rows in reverse order, i.e., Γ_1^R . When 0 and 1 are added respectively to the left we complete Γ_2 . Observe that in order not to obscure the illustration of the building process the 0 has not actually been written in the first two rows. The third block adds four new rows which accommodate the first four rows in reverse order and then the left-most 1 is prefixed. The fourth block illustrates the procedure again to obtain Γ_4 which is given in full in the fourth column. The generalization to 2^N elements is straightforward. This is a recursive algorithmic way to construct the Gray code table.

There are other possibilities for describing an algorithm leading to Gray codes. For example, once we have chosen the number of bits, N , we can start with an arbitrary string of length N . The next term in the sequence is obtained by changing only one bit. From here on the same procedure is repeated with the proviso that strings which have already appeared in the listing are discarded. Observe however that in this way we certainly obtain a binary Gray code, although we cannot ensure it is the native BRGC.

Apart from these recursive and constructive ways to obtain a BRGC of given length there is also an explicit method [2]. Let us write the k th entry of the BRGC of length N as the string $b_1^N(k)b_2^N(k)\cdots b_N^N(k)$, then the j th bit (from left to right) is given by the parity of a binomial coefficient by the following expression which will be helpful in section 5:

$$b_j^N(k) = 2^{N-2^{N-j}-1} C_{[2^N-2^{N-j-1}-(k+1)/2]} \pmod{2} \quad (2)$$

where $[\cdot]$ is the integer part function and $1 \leq j \leq N, 0 \leq k \leq 2^N - 1$. A comment may be in order here. The previous equation could give the wrong impression that the Gray code representative of the natural number k may sensibly change with the total number of bits N .

Either from equation (2) or directly from (1) one can easily see that this is not the case. Increasing N only adds 0s to the left of the representative string.

The native code is cyclic in the sense that 16, in our example, would go onto 0. More generally, the code is cyclic provided we deal with a number of elements that is a power of 2. It is worth mentioning here that the exact number of cyclic Gray codes for general N , which is related to the number of Hamiltonian paths in a hypercube, is still an open problem in combinatoric analysis. From the numeric standpoint, the Gray code listings that we use here were generated with the short Fortran routine in [20].

Eventually, we are particularly interested in interpreting the entries in the Gray code column as natural numbers written in base 2. Their decimal conversion is given in the last column of table 1. A comparison of first and fifth columns admits the following interpretation which we express in two obviously equivalent ways:

- (i) The last column in table 1 is a specific permutation of the 16 elements $\{0, 1, 2, \dots, 15\}$. Or in general, the length N BRGC gives a permutation of the 2^N elements $\{0, 1, 2, \dots, 2^N - 1\}$ which we denote by P_G .
- (ii) The last column in table 1 defines a function of a discrete variable in a definition domain with 16 (in general 2^N) points.

The reason we have stated separately these otherwise interchangeable sentences is because each of them suggests a different line of approach. In fact, while the permutation terminology is more amenable to combinatorial analysis, the functional sense calls to mind the idea of discrete dynamical system associated with the iteration of functions and related fractal structures. Both interpretations will be considered in this paper and yield interesting findings when considering 2^N elements with $N \gg 1$. By now, we stress the fact that all of our development pivots about the correspondence between the first (decimal sequence in standard order) and the fourth columns (decimal equivalent of the binary Gray code entry) in table 1. Gray codification has been then just an intermediary in the process.

We introduce now some notation. Let $\mathcal{D}^{[N]} \equiv \{D_n = n, n = 0, 1, 2, 3, \dots, 2^N - 1\}$ denote the sequence of the first 2^N natural numbers in base 10 and $\mathcal{G}^{[N]} \equiv \{G_n, n = 0, 1, \dots, 2^N - 1\}$, the decimal representation of the Gray code, i.e., for $N = 4$, the last column in table 1. P_G denotes the permutation that connects $\mathcal{D}^{[N]}$ and $\mathcal{G}^{[N]}$ for a given N , i.e., P_G is the element of the symmetric group S_{2^N} which acts on $\mathcal{D}^{[N]}$ according to

$$P_G(n) = G_n, \quad n = 0, 1, \dots, 2^N - 1. \quad (3)$$

In the next two sections we present a series of heuristic considerations on P_G deferring to section 5 its more technical details.

3. A self-similar Gray set

Let us now represent graphically the action of P_G . We consider then the set \mathcal{A}_N of the \mathbb{R}^2 plane consisting of the points with coordinates $(n, P_G(n)) = (D_n, G_n)$, $n = 0, 1, \dots, 2^N - 1$. In figure 1, we plot $\mathcal{G}^{[N]}$ versus $\mathcal{D}^{[N]}$ up to $N = 18$, where we have connected successive points only as a visual help. A direct view of the disconnected set is given in the lower inset of the same figure. When no confusion may arise we shall continue to refer to this set as the Gray code curve. We could describe it, at first, as a kind of crinkly devil's staircase. In view of the algorithmic way explained above to build up $\mathcal{G}^{[N]}$, it becomes clear the self-similar nature of the plot: increasing a new bit to the left-most position of the Gray number means to sum a certain power of 2, and besides, the preceding binary structure is inherited with a reversed ordering. In the \log_2 - \log_2 representation of the upper inset the scale invariance is

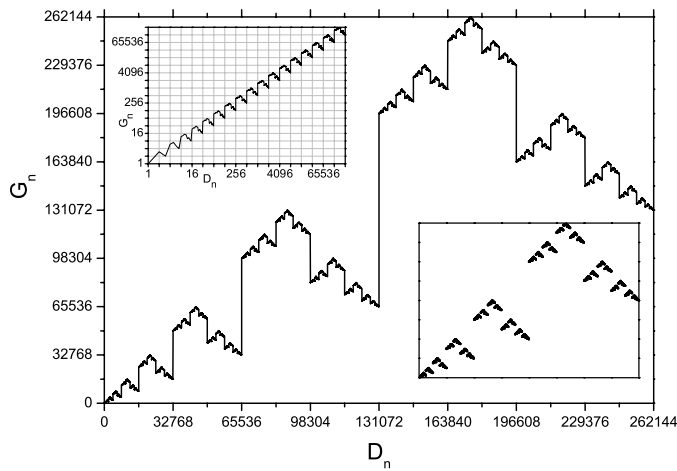


Figure 1. Self-similar Gray code curve.

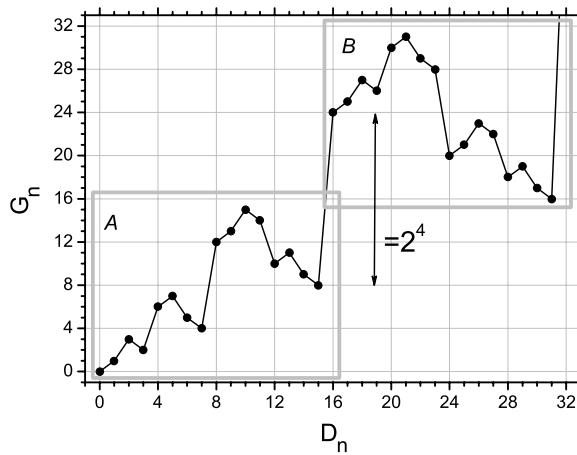


Figure 2. Self-similar construction of the Gray code curve.

fully appreciated. In the present section we describe the set \mathcal{A}_N for large N from the point of view of fractal geometry.

3.1. The Gray code curve as a fractal

The self-similarity of the Gray code curve may be made explicit. Figure 2 contains its first 32 points. Two subsets of 16 points each have been framed. The subset of points in frame B is geometrically obtained from subset A by a procedure of the type *replicate, reflect and translate*, as follows:

- (i) Make an adjacent copy of the set of points in panel A.
- (ii) Reverse the abscissae order in the replicated copy, i.e., substitute it by its image in a vertical mirror.
- (iii) Add 2^4 (=cardinal of the subset) units to every ordinate. This is panel B.

It is easy to see that inside plot frame A a similar procedure has been followed starting with the first eight points. The seed of the procedure was just $\mathfrak{A}_0 = \{(0, 0)\}$.

To put this description in formulae let \mathfrak{A}_k be

$$\mathfrak{A}_k = \{(n, P_G(n)), n = 0, 1, \dots, 2^k - 1\} \subset \mathbb{R}^2, \quad (4)$$

the set of the first 2^k points of our curve, $k = 1, 2, \dots$. Then,

$$\mathfrak{A}_{k+1} = \mathfrak{A}_k \cup F_k(\mathfrak{A}_k) \quad (5)$$

where $F_k(\mathfrak{A}_k)$ is the image set of \mathfrak{A}_k under the function

$$F_k \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} -1 & 0 \\ 0 & 1 \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} + \begin{pmatrix} 2^{k+1} - 1 \\ 2^k \end{pmatrix} \quad (6)$$

defined on \mathbb{R}^2 which corresponds to the previously described manipulations. By this iterative procedure the whole curve in figure 1 grows out. From a purely geometric point of view the previous function defines an affine transformation. But its translation part depends on k , the generation index, a feature to which we shall come back.

Alternatively, we can inscribe the curve in the unit square with vertices $(0, 0)$, $(1, 0)$, $(1, 1)$, $(0, 1)$ by the following iterative procedure. Let us denote by \mathfrak{B}_0 the set $\{(0, 0)\}$, incidentally $\mathfrak{B}_0 = \mathfrak{A}_0$, and define

$$\mathfrak{B}_{k+1} = M(\mathfrak{B}_k) \cup L_k(\mathfrak{B}_k) \quad (7)$$

where $M(\mathfrak{B}_k)$ and $L_k(\mathfrak{B}_k)$ are respectively the image sets of \mathfrak{B}_k under the affine (still k -dependent in the case of L_k) transformations

$$M \begin{pmatrix} x \\ y \end{pmatrix} = \frac{1}{2} \begin{pmatrix} x \\ y \end{pmatrix} \quad (8)$$

$$L_k \begin{pmatrix} x \\ y \end{pmatrix} = \frac{1}{2} \begin{pmatrix} -x \\ y \end{pmatrix} + \frac{1}{2^k} \begin{pmatrix} 2^{k+1} - 1 \\ 2^k \end{pmatrix} \quad (9)$$

which scale down, reflect and translate the curve. From all these considerations the fractal flavour of the whole construction is evident.

3.2. IFS generation of the Gray curve

The procedures just described to construct the Gray set are clearly reminiscent of the so-called iterated functions systems (IFS), very popular in fractal building and image compression [1]. As it is well known a collection of functions are applied, either following a deterministic plan or randomly, to a simple object. The two more striking features of this approach are, on one hand, how astonishingly complex images arise from very simple seeds. On the other hand the perfect coincidence in the final product of deterministic and random processes, although more and more familiar every time thanks in part to the popularization of chaos studies, is always odd. Although not necessarily so, in most of the applications of IFS the functions involved are a few affine transformations with constant parameters.

It is natural then to look for an IFS generator of our set \mathfrak{A}_N . The problem is, however, that one cannot find constant parameters for a bidimensional affine transformation capable of generating by iteration our fractal. The reason being the k dependence mentioned earlier. So we are led to look for more general functions, i.e., nonlinear ones, of the coordinates, such that when iterated produce the desired result. But the k dependence continues to be a nuisance. The only way out is to increase the dimension of the space we are working on from 2 to 3. This is exactly the same idea one takes advantage of in, e.g., classical mechanics to convert a time-dependent Hamiltonian system in an autonomous equivalent.

We then propose to extend \mathbb{R}^2 to \mathbb{R}^3 and take as our IFS the following $\mathbb{R}^3 \rightarrow \mathbb{R}^3$ functions:

$$f(x, y, z) = \left(\frac{x}{2}, \frac{y}{2}, z + 1 \right) \quad (10)$$

$$g(x, y, z) = \left(1 - \frac{x}{2} - 2^{-z}, \frac{1+y}{2}, z + 1 \right). \quad (11)$$

The price one has to pay for getting rid of the k dependence is the higher dimension we work on and the nonlinear character of the function g involved. Needless to say that for representation purposes we only plot the projection on the xy plane. Observe that the action of f and g on the z coordinate prevents the existence of fixed points when each of these functions is iterated on its own. In the xy plane when f is repeatedly iterated the results tend rapidly to the origin. It is only a little bit more complicated to see that iteration of g leads in the limit to the point $(1, 1)$. The combined and competing iteration of the two functions is what produces the final figure. The set $\{f, g\}$ constitutes a *bona fide* IFS. The system is deterministic if f and g are repeatedly applied one after the other. When this is done starting with $(0, 0)$, then the whole set \mathfrak{A}_N is generated. It is a random IFS if f and g are iterated with prescribed probability $\frac{1}{2}$ on an arbitrary initial point in the unit square. \mathfrak{A}_N is obtained as the final attractor of the process.

4. A time series from BRGC

Our proposal now is to interpret G_n as successive values in a time series. Conventional statistics (average values, standard deviations and other indicators calculated from higher moments) are of no particular interest for us here because they depend solely on the numerical values G_n and these are simply the first 2^N natural numbers. The order in which the different values appear does not play any role in the computation of those statistics. The information from the time series obtained in this way could be termed as static. Instead, we are interested in the dynamics represented by the series. The ordering then matters and we need other indicators. With this aim we analyse in the present section the series of increments $\{\Delta_n\}$ with $\Delta_n = G_n - G_{n-1}$, $n = 1, 2, \dots$.

4.1. Statistics on Δ_n

The inset in figure 3 shows a sample of the series of increments $\{\Delta_n\}$. By its very definition the average of the first m values is G_m/m and because it is always positive the Gray curve is globally increasing. The statistics on the set of values $\{|\Delta_n|\}$ reveals that their frequency obeys a power law, with the exception of the very small values (see main graph in figure 3). As a matter of fact, the frequency histogram goes like $\sim 1/|\Delta|$, which means that the probability of encountering a jump $|\Delta| < \delta$ grows as $\log \delta$. One could express this fact by saying that the values G_n are far from equidistant. By contrast, it seems appropriate to mention here that in its binary form, BRGC is characterized by any two consecutive elements being equally spaced by a unit Hamming distance.

4.2. Residence times from Δ_n

An often encountered statistic when analysing time series is the so-called residence time or sojourn time, τ . To deal with it, one has to define a dichotomic property of the system such as crossing a predefined threshold and keep track of the time elapsed between successive crossings. In our case, we focus our attention on changes in the sign of Δ_n . Accordingly, we

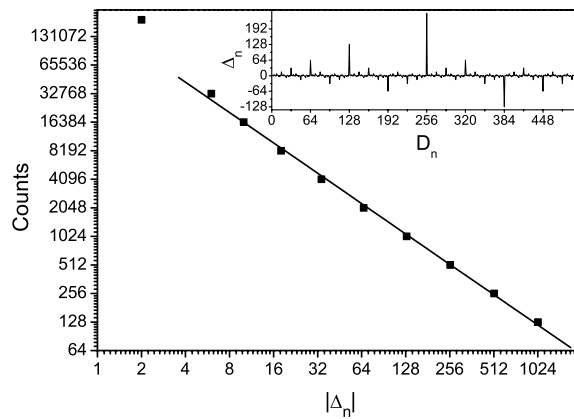


Figure 3. Histogram of the increments $\{|\Delta_n|\}$ in \log_2 - \log_2 scale.

define the residence times as the length of the sequences in which Δ_n keeps its sign. If it stays just one unit in the, say, positive region, we define its sojourn time to be just one unit.

When we collect the statistics for residence times in the set of values of Δ_n the fact which first calls our attention is that τ can only take three different values: 1, 2 or 3. Suppose Δ_n has changed sign with respect to Δ_{n-1} , then there are only three possibilities: either Δ_{n+1} has sign opposite to that of Δ_n (then $\tau = 1$) or Δ_{n+1} and Δ_n have equal sign but different from that of Δ_{n+2} (which corresponds to $\tau = 2$) or Δ_n , Δ_{n+1} and Δ_{n+2} have the same sign which necessarily changes in Δ_{n+3} (giving $\tau = 3$). More than two consecutive steps can never be given without encountering a change of sign. As a matter of fact the statistics reveals more than that: the probabilities (relative frequencies of appearance) for the values of τ are exactly

$$p(1) = 0.25, \quad p(2) = 0.5, \quad p(3) = 0.25, \quad p(\tau > 3) = 0. \quad (12)$$

This behaviour of the residence times contrasts with other known situations. For example, exponentially decreasing probability distributions are common in chaotic systems.

It is worth mentioning that the statistics of residence times we are treating provides information only about changes of sign in Δ_n . It says nothing about the direction (upwards or downwards) of the change. The fact that, as mentioned earlier, the average of increments $\{\Delta_n\}$ is positive suggests the increasing overall character of our curve. By simply looking at it at any scale one corroborates that it is globally increasing, in spite of its ups and downs we perceive in more detailed graphs. In the jargon of time series analysis one says that this reflects extremely persistent dynamics which we analyse next from the Hurst exponent point of view.

4.3. Hurst's R/S analysis of Δ_n

The so-called re-scaled range statistical analysis, R/S analysis, is a method designed to reveal long-run correlations or anti-correlations in a complex process. The idea was first introduced in the analysis of some hydrologic time series, but has also been applied in other studies in climatology, geophysics, econometrics, etc.

Following Feder [5], we define the average increment over a period of k units as

$$\langle \Delta \rangle_k = \frac{1}{k} \sum_{r=1}^k \Delta_r = G_k/k. \quad (13)$$

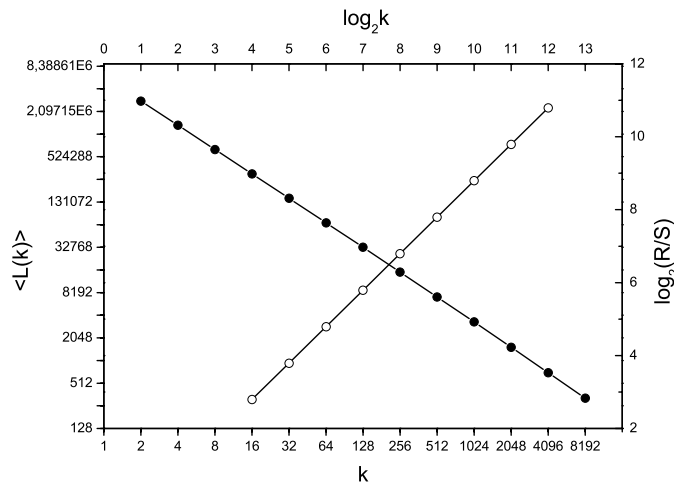


Figure 4. Determination of the fractal dimension (left scale, solid dots) and Hurst exponent (right scale, circles).

The local accumulated departure of Δ_n from the mean $\langle \Delta \rangle_k$ is then

$$X(n, k) \equiv \sum_{r=1}^n [\Delta_r - \langle \Delta \rangle_k] = G_n - \frac{n}{k} G_k. \tag{14}$$

The range of $X(n, k)$ over this period reads

$$R(k) \equiv \max X(n, k) - \min X(n, k) \tag{15}$$

where the extremal functions are evaluated in the interval $1 \leq n \leq k$. The range increases with k . When properly normalized to the standard deviation

$$S = \left(\frac{1}{n} \sum_{r=1}^n [\Delta_r - \langle \Delta \rangle_k]^2 \right)^{1/2} \tag{16}$$

then Hurst found that the re-scaled range R/S is well described by the power law

$$R/S = (k/2)^H \tag{17}$$

where H is called *Hurst exponent*. Purely random series give $H = 1/2$. Curves with $H > 1/2$ ($H < 1/2$) witness a long-run correlation of persistent (anti-persistent) character. This feature means that positive changes are more likely followed by positive (negative) ones, whereas for random phenomena both results for the next move are always equiprobable. Hurst [13] found that most natural phenomena give rise to $H \simeq 0.7$, which is sometimes referred to as Hurst phenomenon.

In the case of the Gray curve a Hurst-like analysis provides a value $H = 0.97(3)$, revealing the extreme persistence of the series of increments, as pointed out above. This analysis appears in figure 4 where R/S (straight line with open circles) is to be read on the right scale.

4.4. Fractal dimension of the Gray curve

We follow a technique by Higuchi [11] to measure the fractal dimension of a set of points $\{n, G_n\}$, $n = 0, 1, 2, \dots, 2^N$, forming the graph of a function like that in figure 1. The idea consists in measuring the average length $\langle L(k) \rangle$ of samples with k points each. Then, if

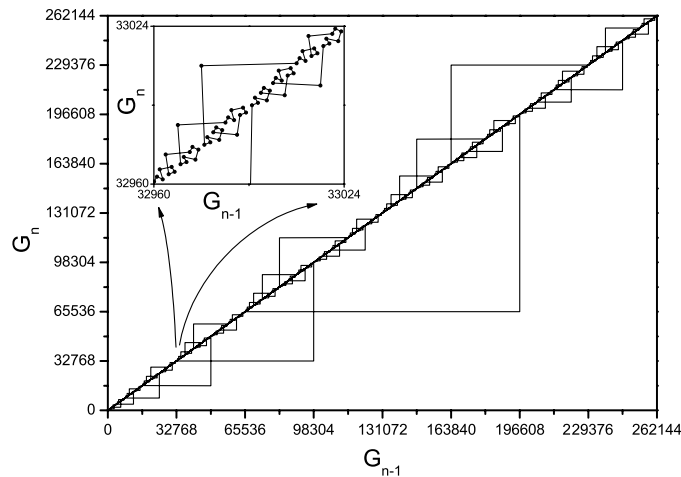


Figure 5. 2D trajectory for the Gray curve and zoom of the marked area.

$\langle L(k) \rangle \propto k^{-d}$, the curve is fractal with dimension d . For instance, a time series obtained with pure noise should give $d = 1.5$, whereas for a regular curve $d = 1$. The method recalls the so-called Richardson plot in fractal geometry.

In figure 4, we have plotted in doubly logarithmic scale $\langle L(k) \rangle$ against k , with $N = 18$ (straight line with solid dots). The straight line is fitted to the points by the least-squares method and gives a slope $d = 1.09(1)$. Here the error estimate comes from the stability of the result with respect to the number of data rather than from the fitting procedure.

These values for H and d are consistent with the connection formula relating H and d which reads $d = 2 - H$.

4.5. Return plot of the Gray curve

A common tool in nonlinear analysis of time series consists in giving a 2D representation of the 1D time series called return plot. A return plot, or first-return map, of the sequence $\{G_n\}$ is a 2D plot of the points $\{(G_{n-1}, G_n)\}$. Thus, for instance, a true random sequence gives rise to a uniformly scattered plot. Chaotic sequences where the generation of the n th term involves only the $(n - 1)$ th term produce 1D patterns. In the present case, the return plot for the Gray curve is shown in figure 5. We have drawn not only the points but also a path connecting them in order to realize the variety of jumps involved, in accordance with the set of values $\{\Delta_n\}$. The inset in figure 5 is a zoom of the indicated zone showing the detail of the curly structure of the curve. It is clear how strongly the natural numbers sequence gets intermingled in the permutation P_G : the recurrence plot for the set $\{D_n\}$ is simply a straight line like $y = x + 1$.

4.6. Flying distances in the random IFS Gray fractal

At the end of subsection 3.2, it was pointed out that the Gray fractal can be obtained as the attractor of a random process of iteration of functions f and g . Here, we investigate the statistics obeyed by the distances travelled in every move. This point of view quits the static perspective on the fractal and enters a new kind of dynamics.

Starting from an arbitrary point in the unit square, runs of 2^{22} points were carried out. We have retained just the last 2^{21} points, and the previous ones were considered as a transient.

The idea is to lose any memory about the initial position by getting close enough to the true attractor. A histogram of the Euclidean distances travelled between two consecutive points in a random walk on the Gray fractal is shown in figure 8. The presence of equally spaced peaks on the log scale is worth noting. The same is true for the bottom of valleys. This phenomenon means that the geometry of the fractal selects (or avoids) some particular flying distances along its random generation. The specific mechanism underlying this selection rule is unclear, as according to the statistic of Δ_n in subsection 4.1 all length distances are present in the fractal itself.

5. The native Gray permutation

In this section, we collect some properties of the permutation P_G which have been instrumental in most of the considerations presented so far. In correspondence with the name *native Gray code* sometimes used for the BRGC we call P_G the *native Gray permutation*. A little experimentation and some reflection lead to the following observations:

- (i) The recurrent definition of the permutation P_G as a function with domain $\mathfrak{D}^{[N]}$ is

$$P_G(2^n + m) = 2^n + P_G(2^n - m - 1) \tag{18}$$

with $P_G(0) = 0$, and where $n = 0, 1, \dots, N - 1, m = 0, 1, \dots, 2^n - 1$. Particularly interesting values are $P_G(2^n - 1) = 2^{n-1}$ and $P_G(2^n) = 3 \times 2^{n-1}, n \geq 1$.

- (ii) An explicit expression for $P_G(n)$ can be deduced from equation (2). Given $n \in \mathbb{N}, n \geq 1$, the minimum number of bits necessary to write it in base 2 is $v(n)$ given by

$$v(n) = [\log_2(n)] + 1.$$

Then, we have

$$P_G(n) = \sum_{j=1}^{v(n)} b_j^{v(n)}(n) 2^{v(n)-j} \tag{19}$$

where $b_j^{v(n)}(n)$ has been defined in section 2

- (iii) The equation $P_G(x) = x$ has, for any value of N , only two solutions: $x = 0$ and $x = 1$. It is to say, 0 and 1 are the only elements of $\mathfrak{D}^{[N]}$ which P_G leaves invariant. All other elements of $\mathfrak{D}^{[N]}$ under the action of P_G are grouped in cycles, but only cycles of length 2^l will appear. In symbols it means that we can factor out P_G in cycles in the form

$$P_G = \prod_{l=0}^{k(N)} c_l^{n_l} \tag{20}$$

where c_l represents a cycle of 2^l elements of which n_l different specimens appear in P_G . $k(N)$ indicates that the length of the longest cycle or cycles is $2^{k(N)}$.

- (iv) The distribution of the 2^N elements of $\mathfrak{D}^{[N]}$ in cycles follows a strict procedure: they fill cycles c_l in increasing order of l . No 2^l -cycle appears until all allowable 2^{l-1} -cycles are filled up. Moreover, n_l has a maximum value independent of N . In other words, once a certain number of 2^l -cycles has been reached no matter how big N becomes no other 2^l -cycle will be formed. Furthermore, when N increases the factorization of P_G keeps the structure corresponding to the lower values of N . The new cycles formed are as long as the longest ones for previous values of N or longer. This property confers to P_G a certain type of shell structure.

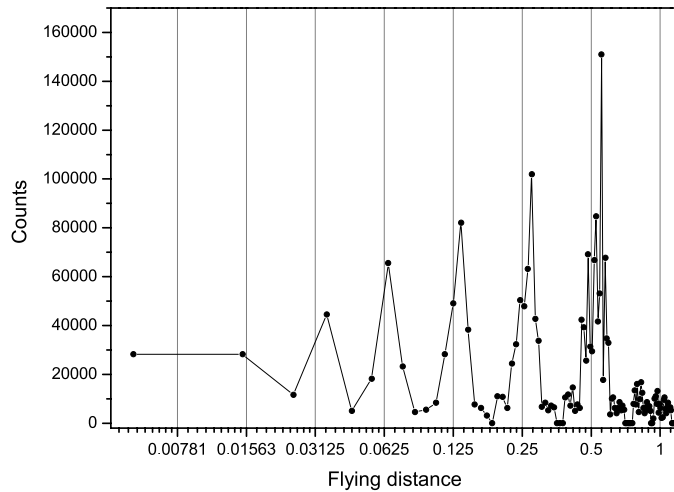


Figure 6. Histogram of flying distances in the random IFS Gray fractal with $N = 22$ and bin size 0.02.

(v) 2^l -cycles will appear for the first time for $N = 2^{l-1} + 1$. In the language of the theory of dynamical systems one could say that at these values of N a period doubling bifurcation appears. This way of speaking has to be considered simply as an analogy and cannot be taken too literally. It is illustrated with an example in figure 6 which is self-explanatory.

With all these considerations in mind a careful bookkeeping shows that the length of the longest cycle in P_G is $2^{k(N)}$ with

$$k(N) = \left\lceil \frac{\ln(N - 1)}{\ln 2} \right\rceil + 1 = [\log_2(N - 1)] + 1, \quad N \geq 2 \tag{21}$$

where $[\cdot]$ stands for integer part. Obviously, for $N = 1$ we have $k(1) = 0$ and, for any N , $n_0 = 2$.

According to observation (iv) above, the maximum value of n_l, n_l^{\max} is, for $l < k(N)$, independent of N and it is not difficult to see that

$$n_l^{\max} = 2^{(2^{l-1}-1)}(2^{(2^{l-1})} - 1), \quad 0 < l < k(N) \tag{22}$$

while the number of longest cycles is

$$n_k = 2^{(2^{k-1}-k)}(2^p - 1) \quad p = N - 2^{k-1} \tag{23}$$

where in these formulae $k(N)$ has been for simplicity written as k . Observe that in this last case the dependence on N is explicit. As a test one can check that

$$\sum_{l=0}^{k(N)} 2^l n_l = 2^N. \tag{24}$$

To fully appreciate the complexity of the cycle structure of the Gray permutation P_G we show in table 2 a few cases of values characterizing that structure. It gives the complete cycle structure of P_G up to $N = 10$ and includes also the total number of cycles, N_c , and the average cycle length, l_{av} .

The reminiscence of the atomic shell model mentioned earlier is of course purely symbolic but it is reinforced when observing figure 7. In it we represent $k(N)$ as a function of N and one can see that the jumps occur for values of N which are exact powers of 2. The inset shows that precisely at these values the number of longest cycles falls down.

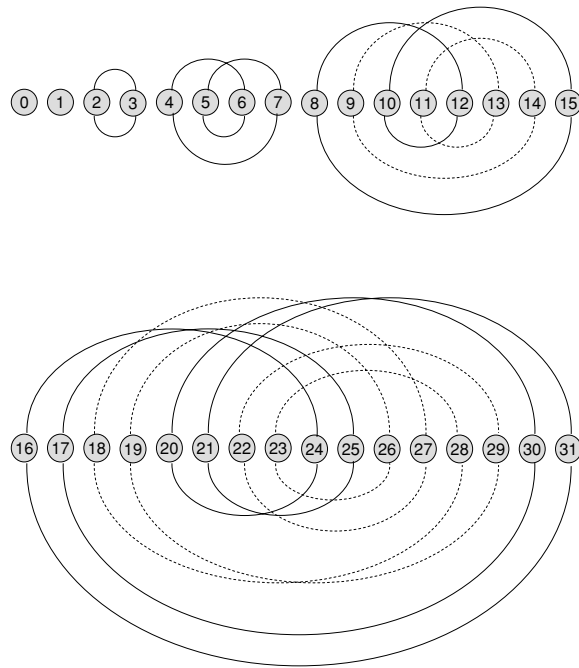


Figure 7. Orbits for Gray permutation of elements 0–31.

Table 2. Cycle structure of permutation P_G .

N	2^N	n_0	n_1	n_2	n_3	n_4	N_c	l_{av}
1	2	2	0	0	0	0	2	1
2	4	2	1	0	0	0	3	1.33
3	8	2	1	1	0	0	4	2
4	16	2	1	3	0	0	6	2.66
5	32	2	1	3	2	0	8	4
6	64	2	1	3	6	0	12	5.33
7	128	2	1	3	14	0	20	6.4
8	256	2	1	3	30	0	36	7.11
9	512	2	1	3	30	16	52	9.85
10	1024	2	1	3	30	48	84	12.19

Of course, the general formulae we have deduced do not give explicitly the elements in each cycle like for example, for $N = 4$,

$$P_G = (0)(1)(2, 3)(4, 6, 5, 7)(8, 12, 10, 15)(9, 13, 11, 14),$$

but this information is much less relevant for our purposes.

It is interesting to note the somewhat distinguished role of P_G in the symmetric group S_{2^N} . It is clear that from the information given in this section one can easily draw the Young diagram corresponding to the equivalence class containing P_G in the group S_{2^N} . It is also obvious that binary Gray codes other than BRGC originate also permutations in S_{2^N} , but these may have a cycle structure very different to that of P_G and consequently behave in other ways. On the other side, any element $Q \in S_{2^N}$ generates trivially a binary code, but in general it is not of the Gray type.

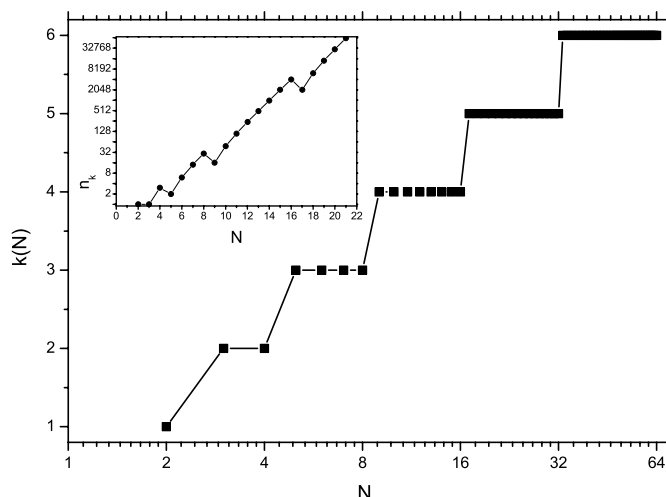


Figure 8. Statistics of cycles as a function of N .

6. Final comments

Among combinatorics problems with physical applications the study of permutations has always been of particular interest. As recent exponents of it, in [6] and [18] random permutations are used in connection with stochastic processes. On the other hand, Gray codes have also met applications in different branches of physics, as widely shown in the introduction.

From the point of view of combinatorics or graph theory relationships between this type of codes and permutations are manifold and are being explored in different directions in the mathematical literature [3, 9, 23]. But from a physics standpoint no such kinship seems to have been studied in detail. In this paper, we have analysed the permutation associated with the binary reflected Gray code from the perspective, and with the techniques, of a dynamical system and its corresponding time series. In so doing we have been led in a natural way to enhance some fractal aspects of the system. One may hope that this line of approach could be helpful in understanding a more dynamical role of permutations.

The results of the analysis carried out in this paper involve only integer arithmetic, some bookkeeping and graphical considerations and so constitute, in our opinion, appropriate material for pedagogical purposes.

As a final summary we would like to mention as a particularly appealing result of our approach the generation of the permutation P_G by an iterated function system applied either in a deterministic or a random way. Also worth mentioning is the discovery of strong selection rules acting on the system which manifest themselves in the distribution of residence times with only three allowed values and of flying distances. Which, if any, of these results apply to other permutations associated with other codes may perhaps be worth considering in the future, as would a deeper understanding of the selection rules mentioned before.

Acknowledgments

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